

# Run II Emittance Preservation

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FNAL

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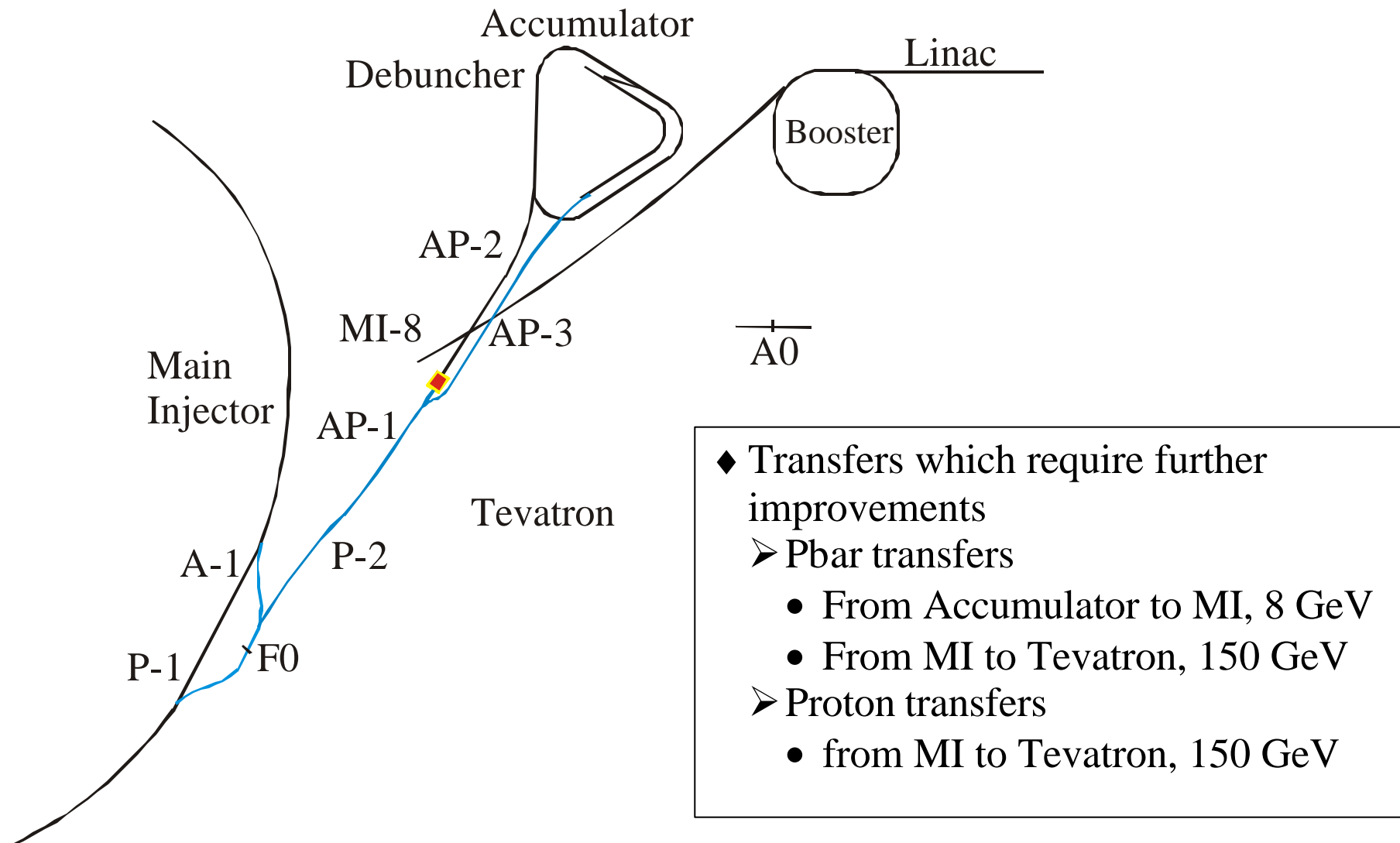
## *Talk outline*

1. Sources of emittance dilution
2. Beam transfers
  - a. Injection errors
  - b. Optics mismatch
3. Luminosity lifetime
- Conclusions

# 1. Sources of the Emittance Dilution

- ◆ **Emittance growth is irrecoverable and all the means need to be applied to prevent it**
  - No antiproton beam cooling after the beam leaves accumulator
  - No cooling for proton beam
- ◆ **Emittance growth reduces initial luminosity and, consequently, integrated luminosity**
  - due to beam size increase
  - due to intensity loss related to scraping particles with large amplitudes
- ◆ **Emittance preservation requires**
  - Prevention emittance growth during beam transfers
    - Prevention/suppression of the emittance growth due to injection errors
      - Turn-by-turn measurements position measurements for every injected bunch are used to close the orbit
      - Injection dampers
    - Careful optics design and measurements to prevent emittance growth due to optics mismatches
  - Prevention emittance growth at beam acceleration
    - Tune and chromaticity control
  - Prevention/Reduction of the emittance growth during store
    - Gas scattering
    - Noises affecting the beam (RF, kickers, *et. c.*)
    - IBS scattering
    - Beam-beam effects

## 2. Beam Transfers



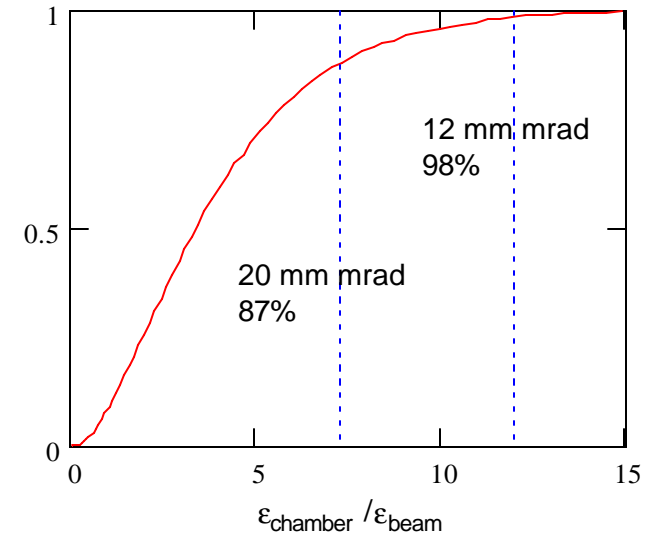
The layout of transfer lines

## Status and projections for the beam transfers

	Line	Energy [GeV]	March 2003		Goal (~June 2003)	
			$\epsilon_{\text{in}} / \epsilon_{\text{fin}}^*$ [mm mrad]	Transfer efficiency	$\epsilon_{\text{in}} / \epsilon_{\text{fin}}^*$ [mm mrad]	Transfer efficiency
Accumulator to MI, $\bar{p}$	AP3-P1	8	6 / 8	97%	7 <sup>**</sup> / 8	>99%
MI to Tevatron, $\bar{p}$	A1	150	10/17	99%	10 / 12	99%
MI to Tevatron, p	P1	150	21/24	99%	21 / 23	99%

\*  $\epsilon \equiv (\mathbf{e}_x + \mathbf{e}_y)/2$ ;    \*\* 170 mA pbar stack;

- ◆ Pick luminosity rises with emittance decrease due to
  - Smaller beam size at collisions
  - Beam current increase due to increase of beam life time at the injection and top energies
    - Reducing pbar emittance in the Tevatron from ~20 mm mrad to ~12 mm mrad should decrease pbar loss at acceleration and squeeze from ~10-12% to 1-2%
- ◆ Presently, major contribution to the antiproton beam emittance growth comes from the beam transfers



## Emittance Growth due to Injection Oscillations

	Present performance		Goal		
	$A_x/A_y$ [mm]	$\Delta\epsilon_x/\Delta\epsilon_y$ [mm mrad]	A, [mm] no damper	A, [mm] with damper	$\Delta\epsilon$ [mm mrad]
Accumulator to MI, $\bar{p}$	1 – 2	0.5 – 2	< 0.7	3	< 0.25
MI to Tevatron, $\bar{p}$	0.5 – 1	1.2–4.7/ 1.9–7.7	< 0.25	1.5	< 0.5
MI to Tevatron, p	0.25–0.5	0.3–1.2/0.48–1.9	< 0.25	1.5	< 0.5

### ◆ Sources of injection oscillations

- Initial injection errors
- Shot-to-shot field variations in dipoles, dipole correctors and kickers
- Bunch-to-bunch field variations in correctors for both pbar transfers

### ◆ Ways to implement good quality of transfers

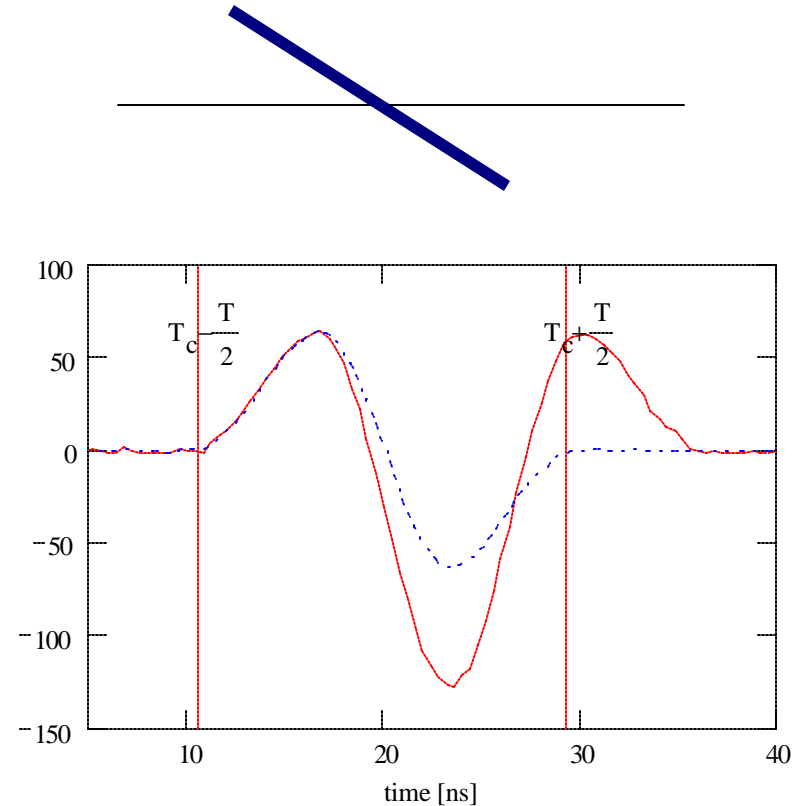
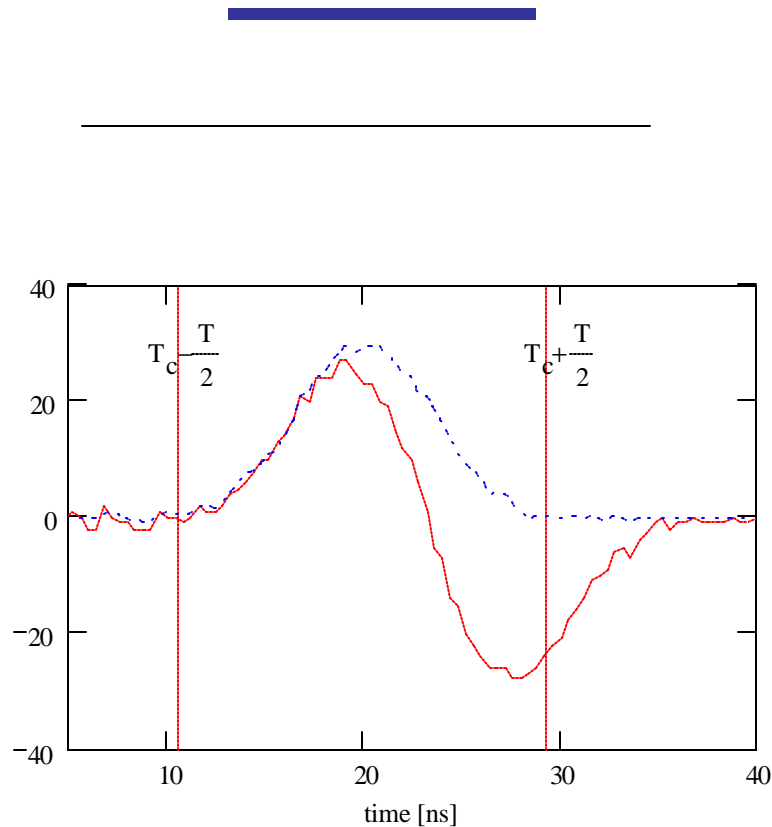
- Orbit closure before colliding bunches are injected
- Orbit closure correction for every new injection
  - Tevatron injection dampers should come in April-May 2002

### ◆ Turn-by-turn BPM measurement for every injected bunch confirms transfer quality

- 3 types of Beam Line Tuners (BLT) were tested before we come up to the final choice
- Old Run I BLT
- BLT based on a digital receiver
- **BLT based on the fast (0.4 ns/sample) digital scope – is a final choice**

◆ Complications of BLT(BPM) turn-by-turn signal analysis come from interplay of

- effects of chromaticity
- and close lengths of the bunch and BLT plates



- Numerical deconvolution of digitized BLT signals create simple and reliable way to compute bunch center of gravity
- In addition to the standard beam position the digital scope data represent internal motion in the bunch

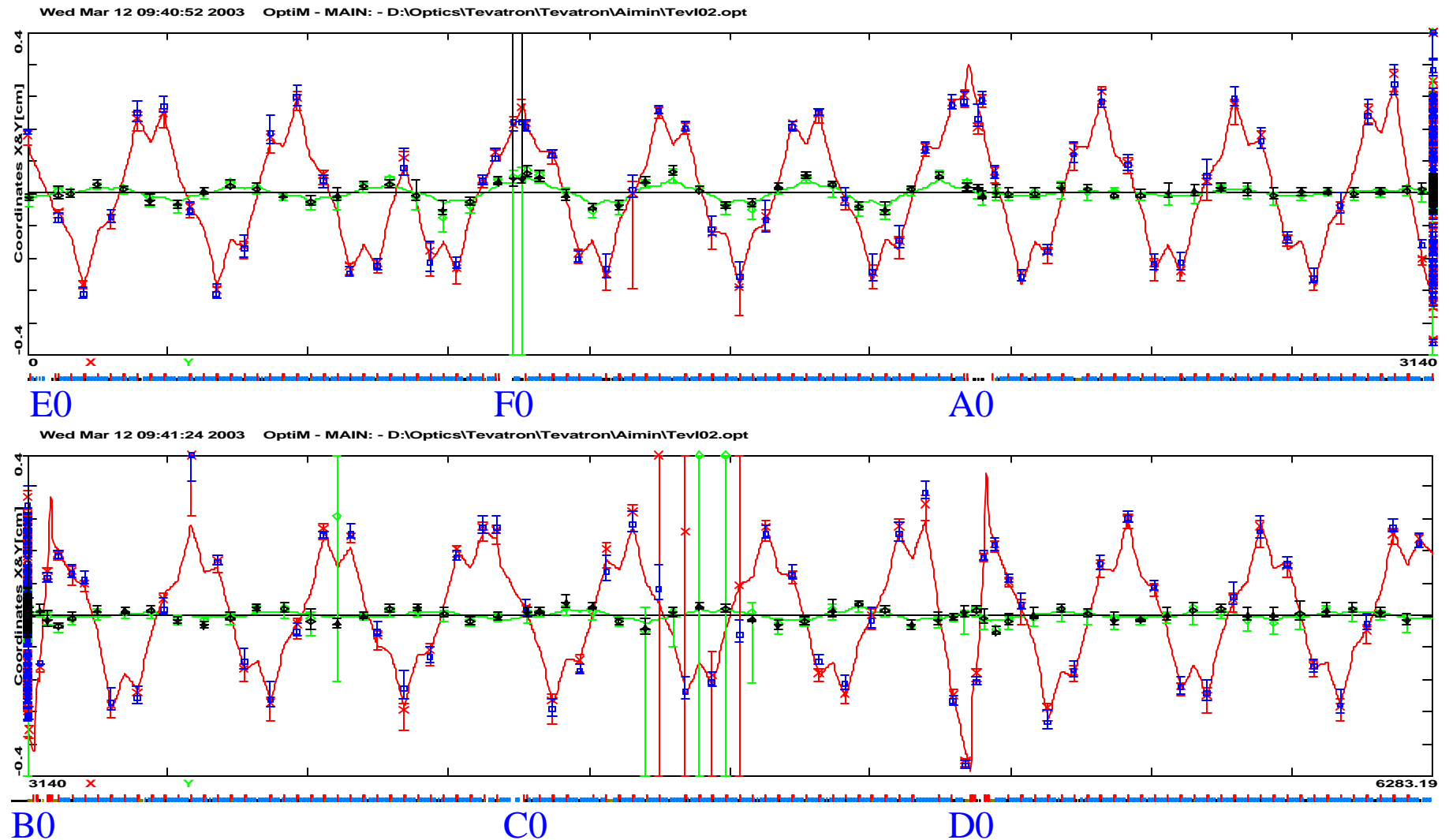
## Optics Mismatch

- ◆ Types of possible optics mismatch
  - Betatron functions mismatch
  - Dispersion mismatch
  - Mismatch due to coupling
  - Optics changes for antiprotons due to long-range beam-beam effects
- ◆ Pbar transfers from MI to Tevatron exhibit significant emittance growth
  - Round trip emittance measurements (MI→Tevatron→MI) are the most reliable.  
They yield
    - $\Delta\epsilon_x \sim 5\text{-}8 \text{ mm mrad}$ , horizontal
    - $\Delta\epsilon_y \sim 3\text{-}5 \text{ mm mrad}$ , vertical
  - All indirect indications point to the A1 line transfer as a major problem
  - The reason of the emittance growth is still not 100% clear
    - Optics for A1 line was corrected, measured and is believed not to be a problem
    - Initial injection oscillations contribute only fraction of the measured increase
    - Tevatron optics and coupling is still not fully understood
      - Optics measurements at central orbit and their analysis were carried out
        - A number of optics problems has been discovered
        - but the difference with the design model is sufficiently small and cannot explain the observed emittance growth
      - Optics measurements at pbar helix will be acquired within a week
        - It is expected that optics at pbar helix is quite different from the central orbit optics and this is a major source of the problems

# Tevatron optics measurements

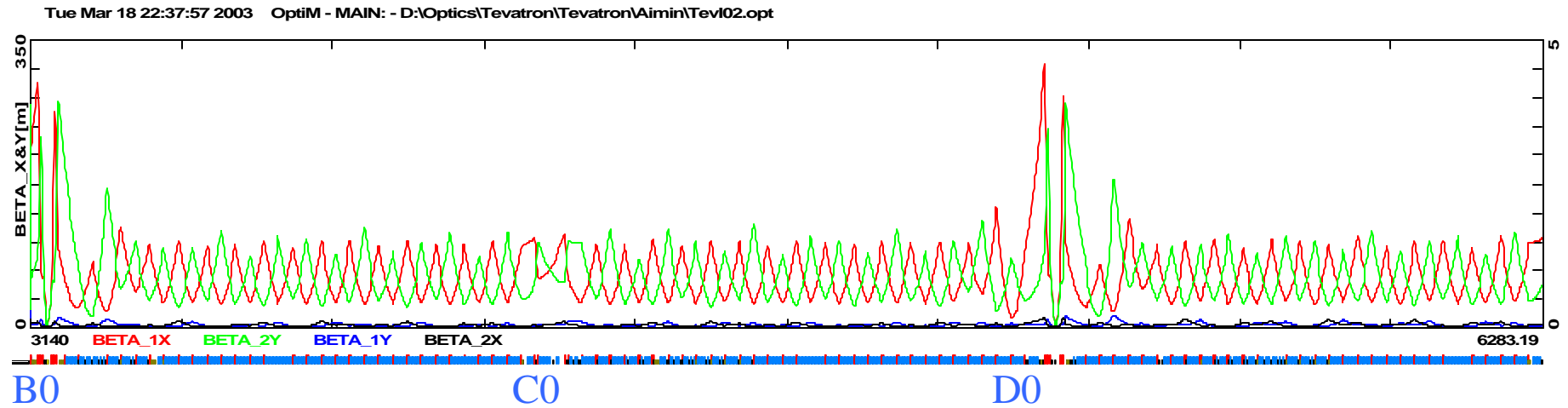
150 GeV, central orbit, data were taken at Feb.20. 2003

X1: HE42 = 50 mrad

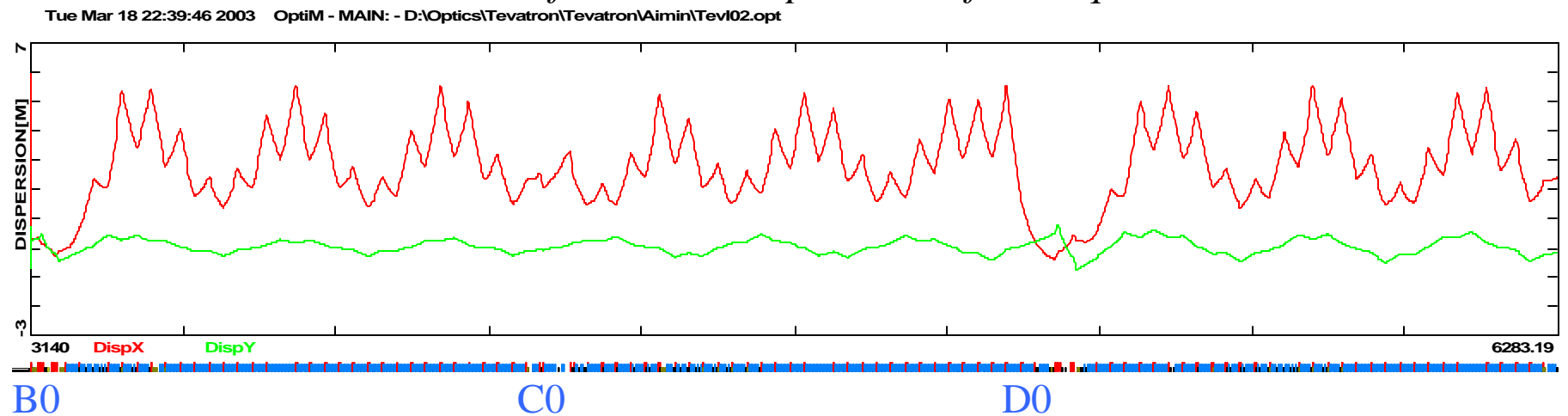


- Altogether five differential orbits are acquired (4 correctors and energy change)
- Fitting to measurements yields reliable model of the machine





*4D beta-functions and positions of skew-quads*



*Dispersions*

Skew-quad term in dipoles ( $\Delta B/B \approx 5.7 \cdot 10^{-5}$  at 1 cm) is the main source of Tevatron coupling

- It originates from the drowning of the SC coil relative to the iron core

## Fudge factors and rolls to fix linear optics

### Global corrections

- $F_{\text{bendq}} = 2\%$ ; correction of dipole edge focusing
- $F_{\text{mq}} = 0.165\%$ ; correction of main bus quad focusing
- $F_{\text{Dskew}} = 1.44$  units; skew quadrupole field of main dipoles

### Point like corrections of quadrupole focusing

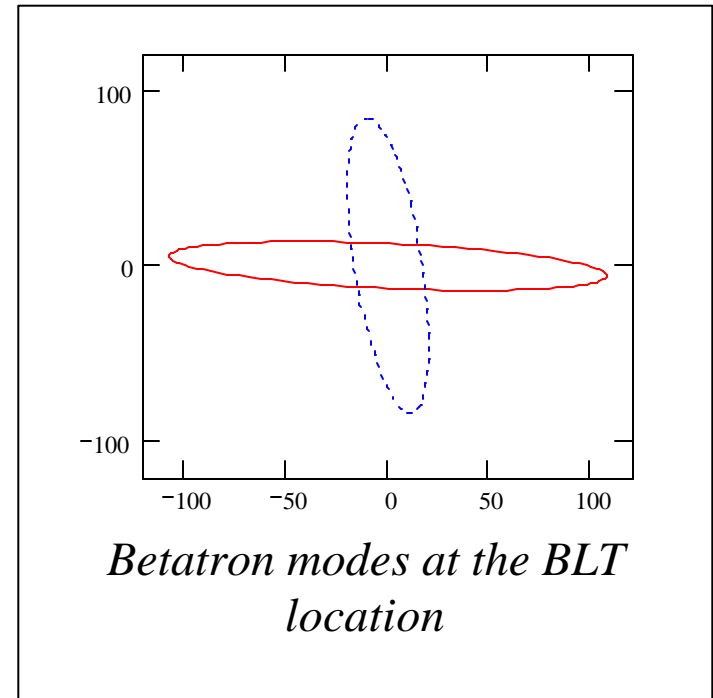
- $F_{\text{qA0U}} = 1\%$ ; related to beam displacement in A0
- $F_{\text{qC27}} = -2\%$
- $F_{\text{CQ7}} = 20\%$ ; that corresponds to 4.4% correction for regular main bus quad
- $F_{\text{B0Q3F}} = 0.37\%$
- $F_{\text{D0Q3F}} = 0.6\%$ ;
- $F_{\text{D0Q2D}} = 1\%$ ;

### Quad rolls

- $Q_{\text{roll\_A0U}} = 0.5$  deg; related to beam displacement in A0
- $Q_{\text{roll\_B0Q7}} = -4$  deg;

## Conclusions for BPMs

- T:VPF0LU and T:VPF0LD are swapped
- T:HPF0LU and T:HPF0LD are swapped
- T:HPC28 has wrong polarity
- T:HPB22 has large noise and incorrect differential position
- T:VPC21 , T:HPC22 and T:HPC36 have large difference for positive and negative bumps. Probably there is large beam offset in BPM.



### 3. Luminosity Lifetime

The model takes into account the major beam heating and particle loss mechanisms

- Phenomena taken into account

- ⇒ Interaction with residual gas

- ◆ Emittance growth due to electromagnetic scattering

- ◆ Particle loss due to nuclear and electromagnetic interaction

- ⇒ Particle interaction in IPs (proportional to the luminosity)

- ◆ Emittance growth due to electromagnetic scattering

- ◆ Particle loss due to nuclear and electromagnetic interaction

- ⇒ IBS

- ◆ Energy spread growth and emittance growth due to multiple scattering

- ⇒ Bunch lengthening due to RF noise

- ⇒ Particle loss from the bucket due to heating of longitudinal degree of freedom

- Phenomena ignored in the model

- ⇒ Beam-beam effects

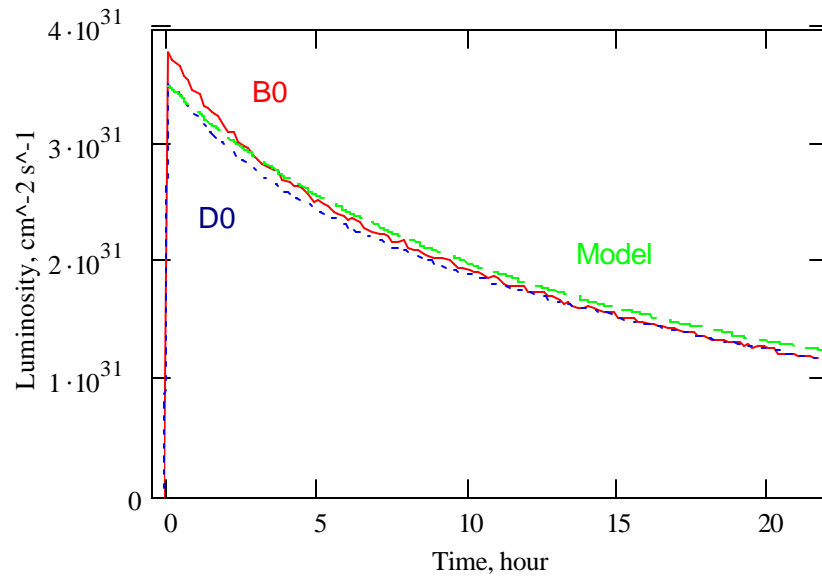
- ⇒ Non-linearity of the lattice

- ⇒ Diffusion amplification by coherent effects

- Thus, it can be considered as **the best-case scenario**

- ⇒ It describes well our best present stores

# Comparison of the Model Predictions to the Store 1953 parameters



$$\epsilon_{px} \cdot 10000 = 19 \text{ mm mrad}$$

$$\epsilon_{py} \cdot 10000 = 19 \text{ mm mrad}$$

$$\epsilon_{ax} \cdot 10000 = 18 \text{ mm mrad}$$

$$\epsilon_{ay} \cdot 10000 = 18 \text{ mm mrad}$$

$$\kappa = 0.3$$

$$\frac{\tau_{\text{gas}}}{3600} = 300.087 \text{ hour}$$

$$d\epsilon/dt_{\text{gas}} = 0.168 \text{ mm mrad/hour}$$

$$\sqrt{d\epsilon^2/dt_{\text{RF}} 3600} = 0.016 \text{ rad/hour}^{1/2}$$

$$N_p = 1.6 \times 10^{11}$$

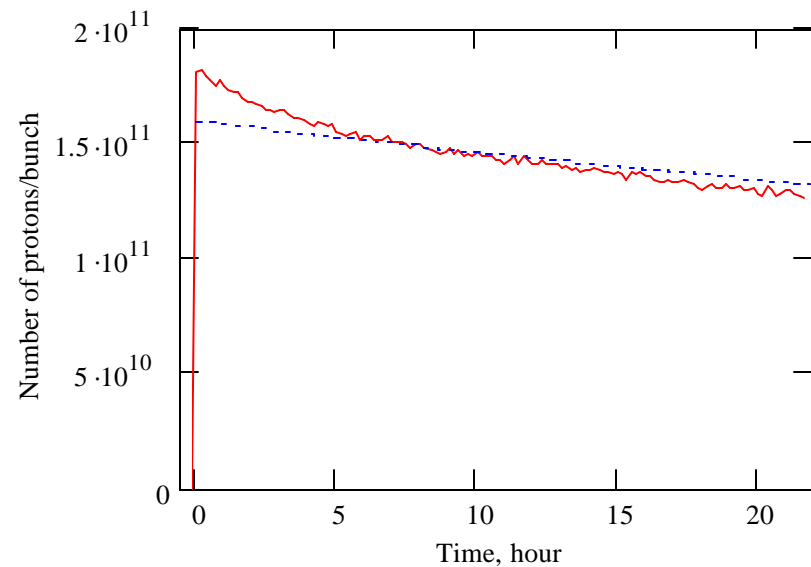
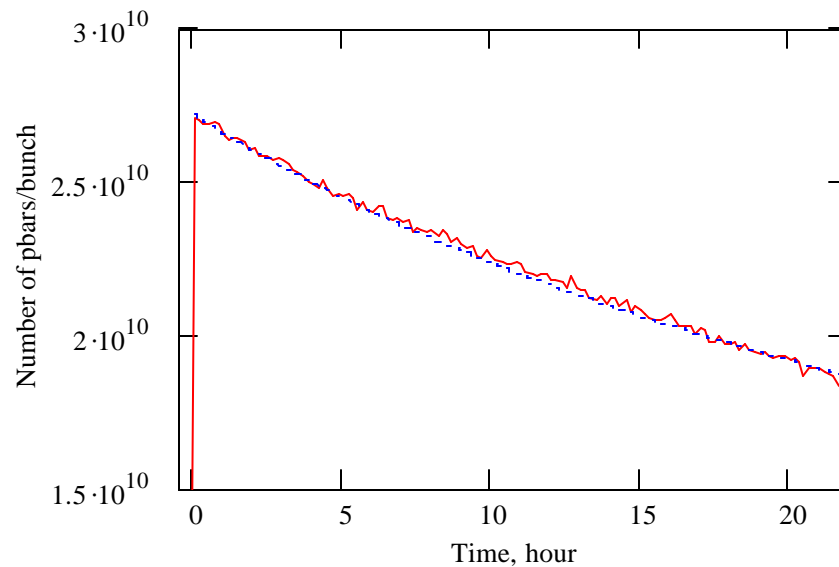
$$N_a = 2.72 \times 10^{10}$$

$$\sigma_s(\sigma_{pp}) = 59.4 \text{ cm}$$

$$\sigma_s(\sigma_{pa}) = 55.112 \text{ cm}$$

$$\text{Lum}_0 = 3.513 \times 10^{31}$$

$$\tau_{\text{Lum}_1} = 13.826 \text{ hour}$$



## Conclusions

- ◆ **Injections errors** have been a leading reason of the antiproton beam emittance growth
  - Introducing turn-by-turn measurements for every bunch injected into Tevatron allowed us to improve orbit closure. That yielded significantly improvement for beam transfers
  - Further improvements are expected after commissioning of Tevatron **injection damper** (April 2003)
- ◆ First **optics correction** in A1 line brought better transfers and luminosity increase
  - Optics corrections and Tevatron, A1 and P1 lines are expected to produce further reduction of the beam emittances
- ◆ Significant progress in understanding of **luminosity lifetime** has been achieved
  - The model predicts that
    - IBS is major mechanism for emittance dilution and beam current reduction
    - There is no other than IBS and major heating mechanisms limiting Tevatron performance
      - Further improvements of Tevatron vacuum and RF noise will not yield significant improvements in integrated luminosity

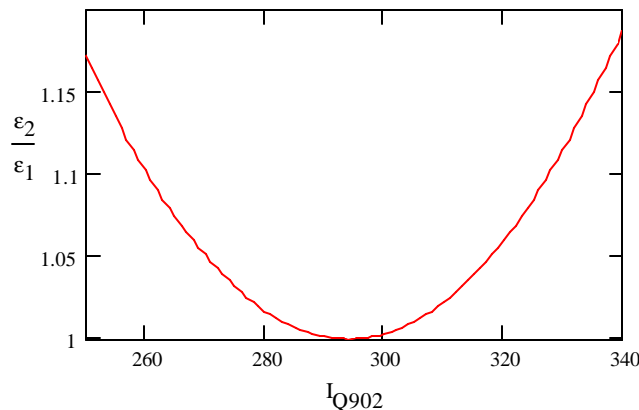
# Emittance Growth due to Betatron and Dispersion Mismatch

Emittance growth from a lattice with  $\mathbf{b}_1$ ,  $\mathbf{a}_1$ ,  $D_1$  and  $D_1'$  to a lattice with  $\mathbf{b}_2$ ,  $\mathbf{a}_2$ ,  $D_2$  and  $D_2'$  is

$$\mathbf{e}' = \frac{\mathbf{e}}{2} \left( \frac{\mathbf{b}_1}{\mathbf{b}_2} [1 + \mathbf{a}_2^2] + \frac{\mathbf{b}_2}{\mathbf{b}_1} [1 + \mathbf{a}_1^2] - 2\mathbf{a}_1\mathbf{a}_2 \right) + \frac{\mathbf{s}_p}{2} \left( \mathbf{b}_2(D_2' - D_1')^2 + 2\mathbf{a}_2(D_2' - D_1')(D_2 - D_1) + \frac{(D_2 - D_1)^2}{\mathbf{b}_2} (1 + \mathbf{a}_2^2) \right)$$

## Emittance growth due single quad focusing error at zero dispersion

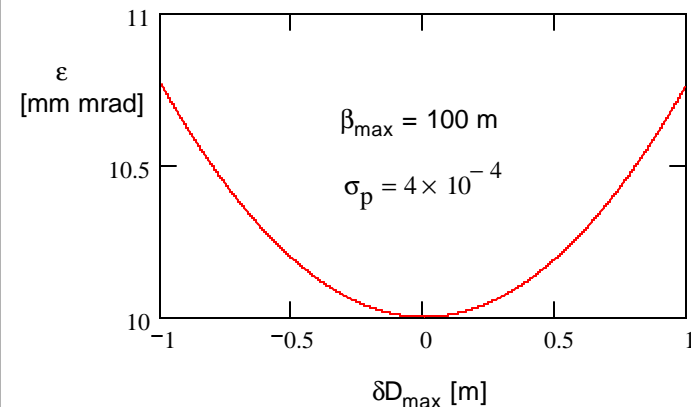
$$\mathbf{e}_2 \approx \mathbf{e}_1 \left( 1 + \frac{d\mathbf{a}^2}{2} \right) \approx \mathbf{e}_1 \left( 1 + \frac{(\mathbf{b}dF)^2}{2F^4} \right)$$



- ◆ Differential orbit measurements allow seeing focusing errors of 1-2%.
  - It is sufficient to tune the line focusing so that the emittance growth would be below 10%.
  - Further improvement is expected from online tuning with orthogonal quads.

## Requirements for dispersion mismatch for MI to Tevatron transfer

$$\mathbf{e}_2 \approx \mathbf{e}_1 \left( 1 + \frac{(\mathbf{s}_p dD_{\max})^2}{2\mathbf{b}_{\max}} \right)$$



- ◆ Dispersion mismatch below about 0.5 m does not produce significant emittance growth

## Emittance Growth due to X-Y Coupling

Emittance growth for beam transfer from an uncoupled lattice with  $\mathbf{b}_x$ ,  $\mathbf{a}_x$ ,  $\mathbf{b}_y$  and  $\mathbf{a}_y$ , to a coupled lattice described by  $\mathbf{b}_{1x}$ ,  $\mathbf{a}_{1x}$ ,  $\mathbf{b}_{1y}$ ,  $\mathbf{a}_{1y}$ ,  $\mathbf{b}_{2x}$ ,  $\mathbf{a}_{2x}$ ,  $\mathbf{b}_{2y}$  and  $\mathbf{a}_{2y}$  with the eigen-vectors

$$\mathbf{v}_1 = \begin{bmatrix} \sqrt{\mathbf{b}_{1x}} \\ -\frac{i(1-u) + \mathbf{a}_{1x}}{\sqrt{\mathbf{b}_{1x}}} \\ \sqrt{\mathbf{b}_{1y}} e^{in_1} \\ -\frac{i u + \mathbf{a}_{1y}}{\sqrt{\mathbf{b}_{1y}}} e^{in_1} \end{bmatrix}, \quad \mathbf{v}_2 = \begin{bmatrix} \sqrt{\mathbf{b}_{2x}} e^{in_2} \\ -\frac{i u + \mathbf{a}_{2x}}{\sqrt{\mathbf{b}_{2x}}} e^{in_2} \\ \sqrt{\mathbf{b}_{2y}} \\ -\frac{i(1-u) + \mathbf{a}_{2y}}{\sqrt{\mathbf{b}_{2y}}} \end{bmatrix}$$

is determined by the following equations:

$$\mathbf{e}_1' = \mathbf{e}_1 A_{11} + \mathbf{e}_2 A_{12}$$

$$\mathbf{e}_2' = \mathbf{e}_1 A_{21} + \mathbf{e}_2 A_{22}$$

$$A_{11} = \frac{1}{2} \left( \frac{\mathbf{b}_x}{\mathbf{b}_{1x}} [(1-u)^2 + \mathbf{a}_{1x}^2] + \frac{\mathbf{b}_{1x}}{\mathbf{b}_x} [1 + \mathbf{a}_x^2] - 2\mathbf{a}_{1x}\mathbf{a}_x \right), \quad A_{12} = \frac{1}{2} \left( \frac{\mathbf{b}_y}{\mathbf{b}_{1y}} [u^2 + \mathbf{a}_{1y}^2] + \frac{\mathbf{b}_{1y}}{\mathbf{b}_y} [1 + \mathbf{a}_y^2] - 2\mathbf{a}_{1y}\mathbf{a}_y \right)$$

$$A_{21} = \frac{1}{2} \left( \frac{\mathbf{b}_x}{\mathbf{b}_{2x}} [u^2 + \mathbf{a}_{2x}^2] + \frac{\mathbf{b}_{2x}}{\mathbf{b}_x} [1 + \mathbf{a}_x^2] - 2\mathbf{a}_{2x}\mathbf{a}_x \right), \quad A_{22} = \frac{1}{2} \left( \frac{\mathbf{b}_y}{\mathbf{b}_{2y}} [(1-u)^2 + \mathbf{a}_{2y}^2] + \frac{\mathbf{b}_{2y}}{\mathbf{b}_y} [1 + \mathbf{a}_y^2] - 2\mathbf{a}_{2y}\mathbf{a}_y \right)$$